

The square potential well

Exact solution and Variational solution using Plane Waves and Gaussians

Report by: Anders Ebro Christensen

Instructor: Hannes Jónsson

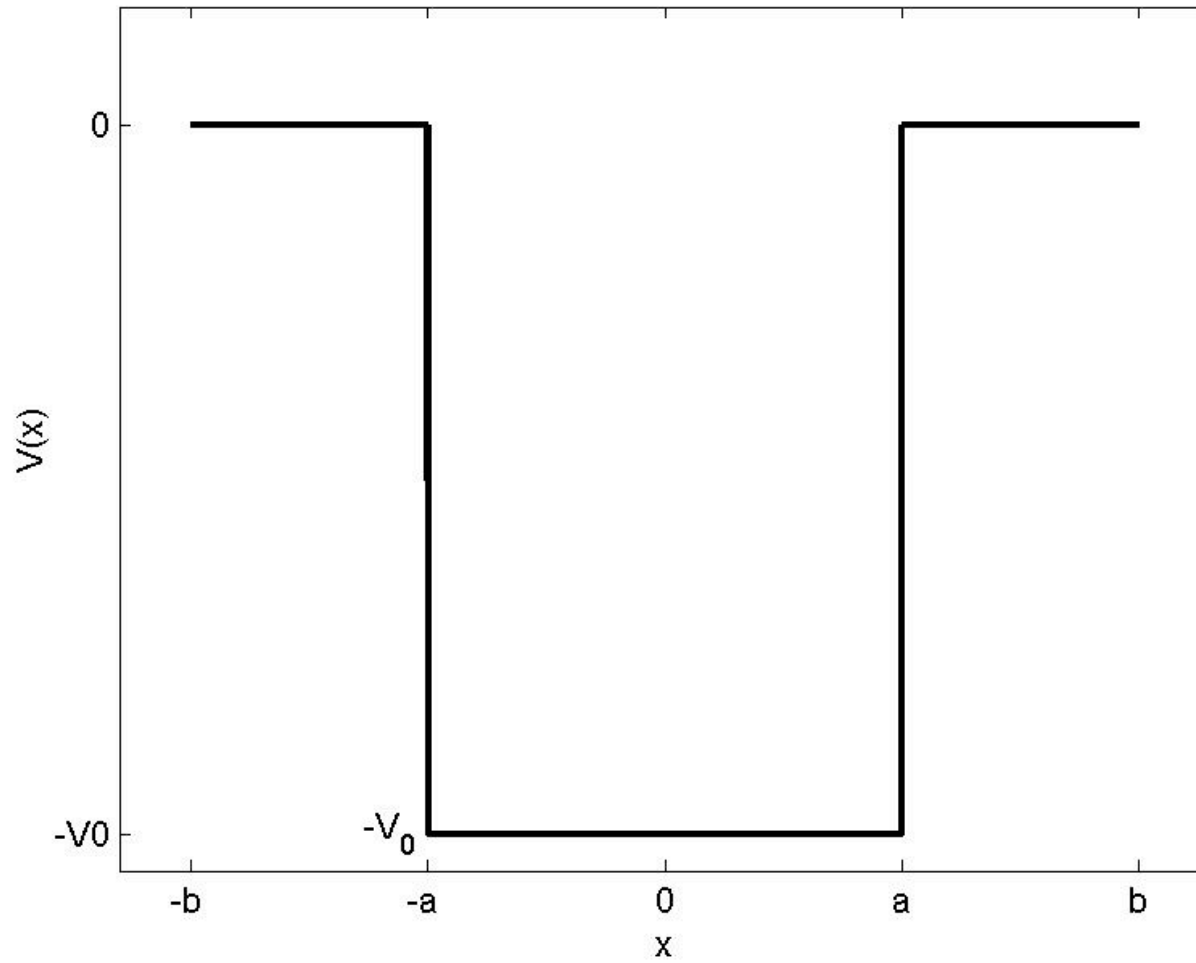
University of Iceland, Department of Chemistry, 2004

Introduction

- We wish to illustrate the methods of density functional theory by going through an example in 1D, and these are the starting steps of the project.
- First Anders will present the square well and the exact solution
- Then Edda will present an approximation based on a basis set of gaussians
- Finally Egill will present an approximation using plane wave as a basis set and compare the results of these different approaches

Square Well

The square well



$$V(x) = \begin{cases} -V_0 & , |x| < a \\ 0 & , |x| > a \end{cases}$$

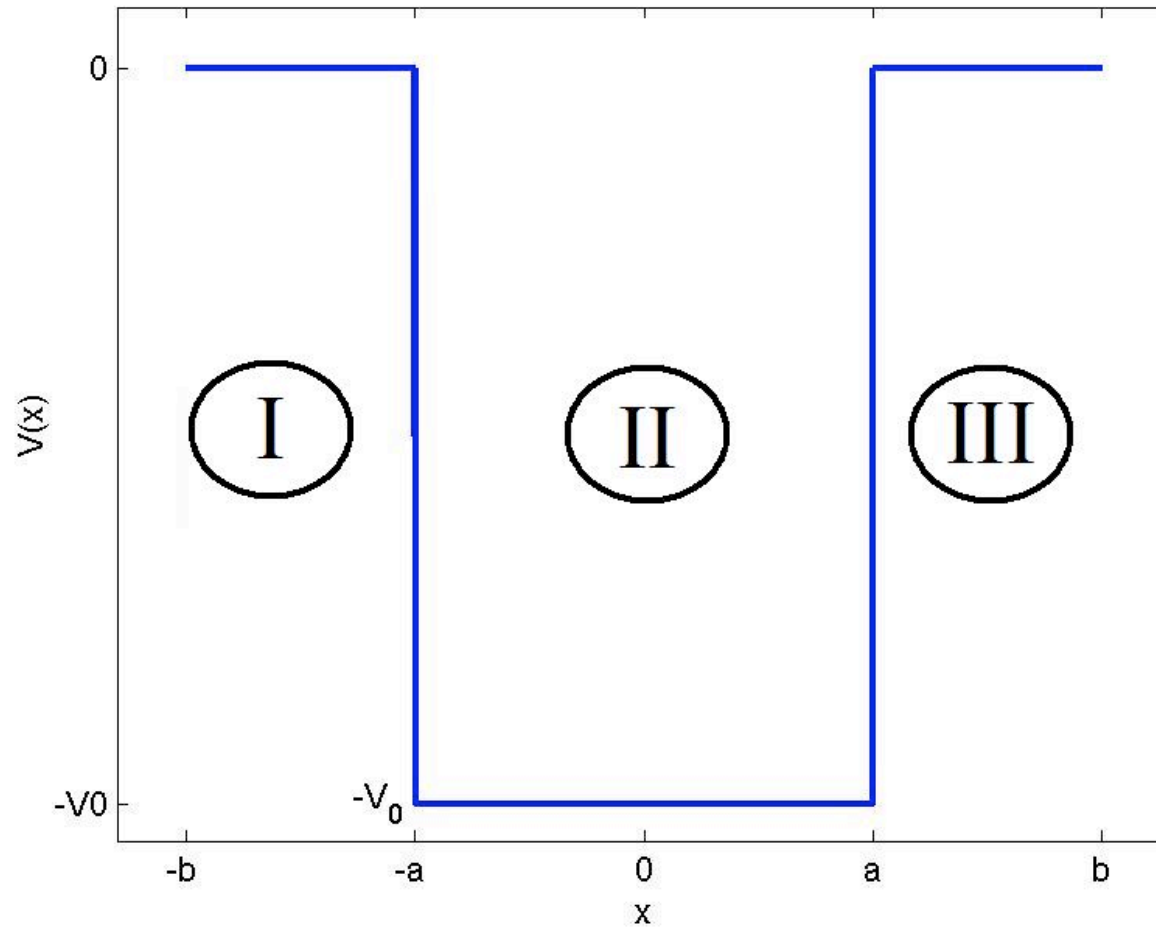
Case: $E < 0$

- In our approximation of the Lithium atom we are only interested in the bound states, and therefore only consider $E < 0$
- The time independent Schrödinger equation is:

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V(x) \right] \psi(x) = E\psi(x)$$

Regions

The Square Well



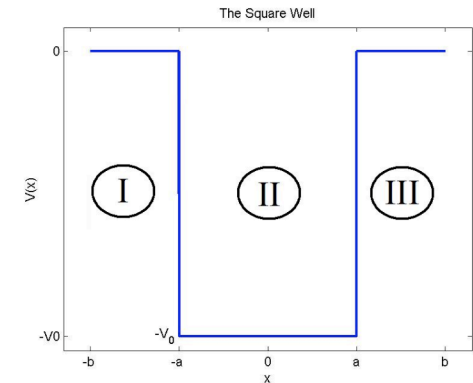
Region I and III

$$a < |x|$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 \right] \psi(x) = E\psi(x)$$

$$\frac{d^2\psi(x)}{dx^2} = -\frac{2mE}{\hbar^2}\psi(x) = \beta^2\psi(x)$$

$$\beta = \sqrt{-\frac{2mE}{\hbar^2}} = \sqrt{\frac{2m|E|}{\hbar^2}}$$



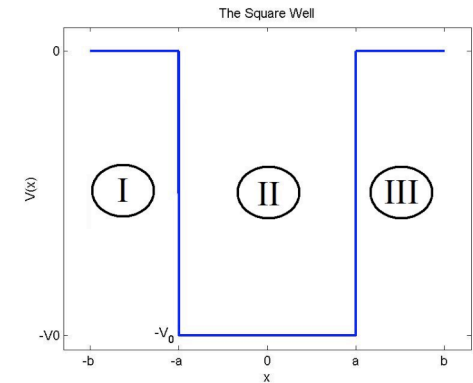
Region II

$$|x| < a$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 - V_0 \right] \psi(x) = E \psi(x)$$

$$\frac{d^2 \psi(x)}{dx^2} = \frac{2m(E + V_0)}{\hbar^2} \psi(x) = \alpha^2 \psi(x)$$

$$\alpha = \sqrt{\frac{2m(V_0 - |E|)}{\hbar^2}}$$



The even solutions

$$\psi(x) = A \cos(\alpha x) \quad 0 < x < a$$

$$\psi(x) = C e^{-\beta x} \quad x > a$$

- By matching up the wavefunctions at the boundaries $x = \pm a$ gives:

$$A \cos(\alpha a) = C e^{-\beta a}$$

$$-\alpha \sin(\alpha a) = -\beta C e^{-\beta a}$$

- Which leads to the equation:

$$\alpha \tan(\alpha a) = \beta$$

The odd solutions

$$\psi(x) = B \sin(\alpha x) \quad 0 < x < a$$

$$\psi(x) = C e^{-\beta x} \quad x > a$$

- By matching up the wavefunctions at the boundaries $x = \pm a$ gives:

$$A \cos(\alpha a) = C e^{-\beta a}$$

$$-\alpha \sin(\alpha a) = -\beta C e^{-\beta a}$$

- Which leads to the equation:

$$\alpha \cot(\alpha a) = -\beta$$

Energy Levels

- The energy levels of the bound states are found by solving the transcendental equations, either numerically or graphically
- We introduce the dimensionless quantities:

$$\xi = \alpha a \quad \text{and} \quad \eta = \beta a$$

- Giving the equations

$$\xi \tan \xi = \eta \quad (\text{for even states})$$

$$\xi \cot \xi = -\eta \quad (\text{for odd states})$$

Energy Levels

- Note that both ξ and η must be positive and such that:

$$\xi^2 + \eta^2 = \gamma^2$$

where

$$\gamma = \sqrt{\frac{2mV_0a^2}{\hbar^2}}$$

γ is often referred to as the strength parameter

Numerical solution

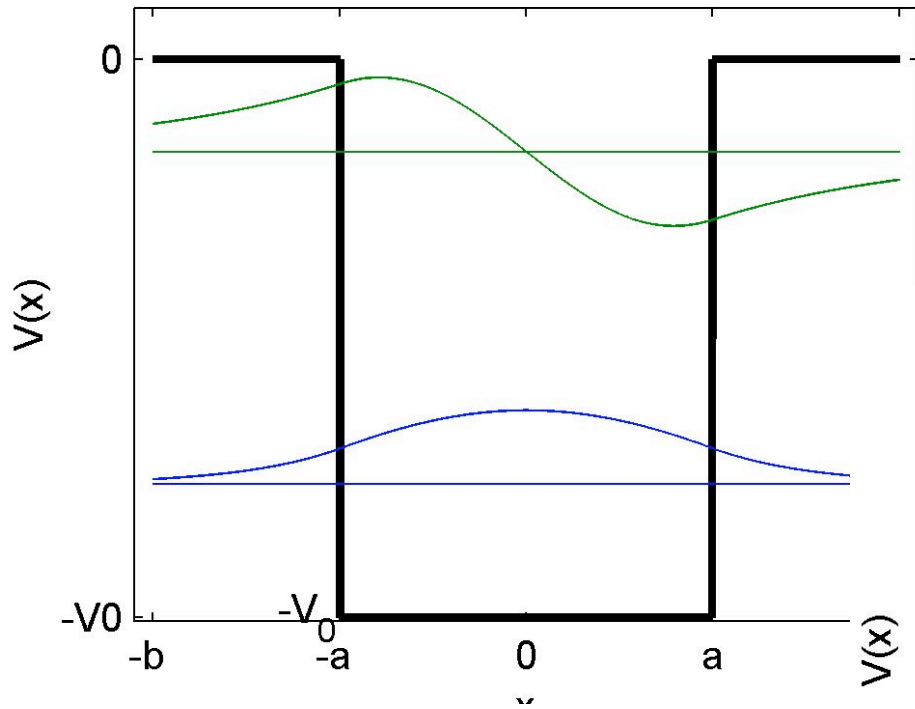
- The equation $\xi^2 + \eta^2 = \gamma^2$ is solved numerically

Results

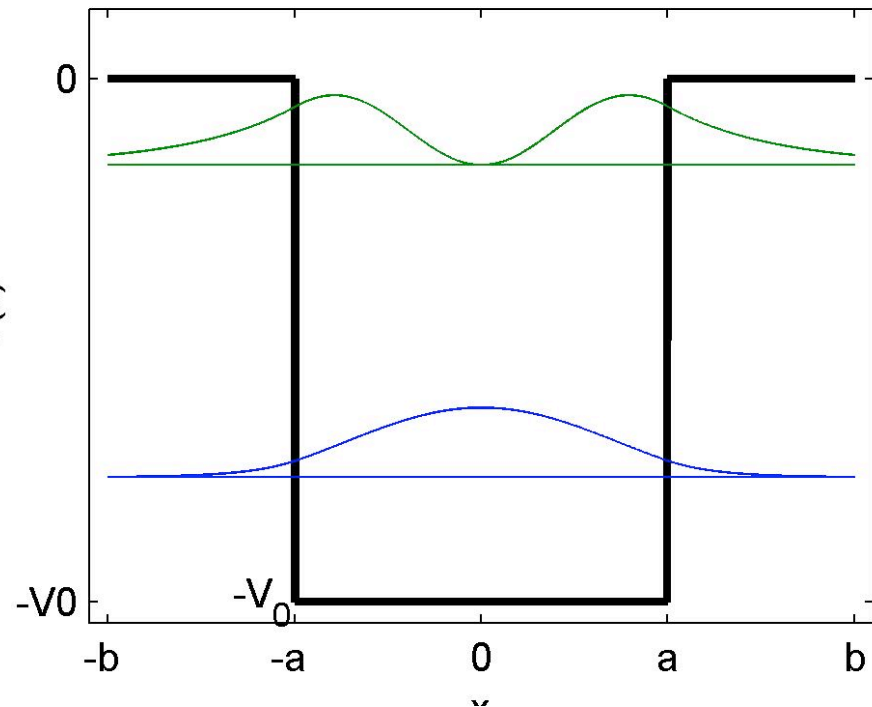
- Using $a=1.55\text{\AA}$ and that the ground state energy should be equal to the ionization energy -5.7eV I get by trial and error that $V_0=7.51\text{eV}$.

Bound states:

The square well with bound states



Probability Density



Strength parameter

$$\gamma = \sqrt{\frac{2mV_0a^2}{\hbar^2}}$$

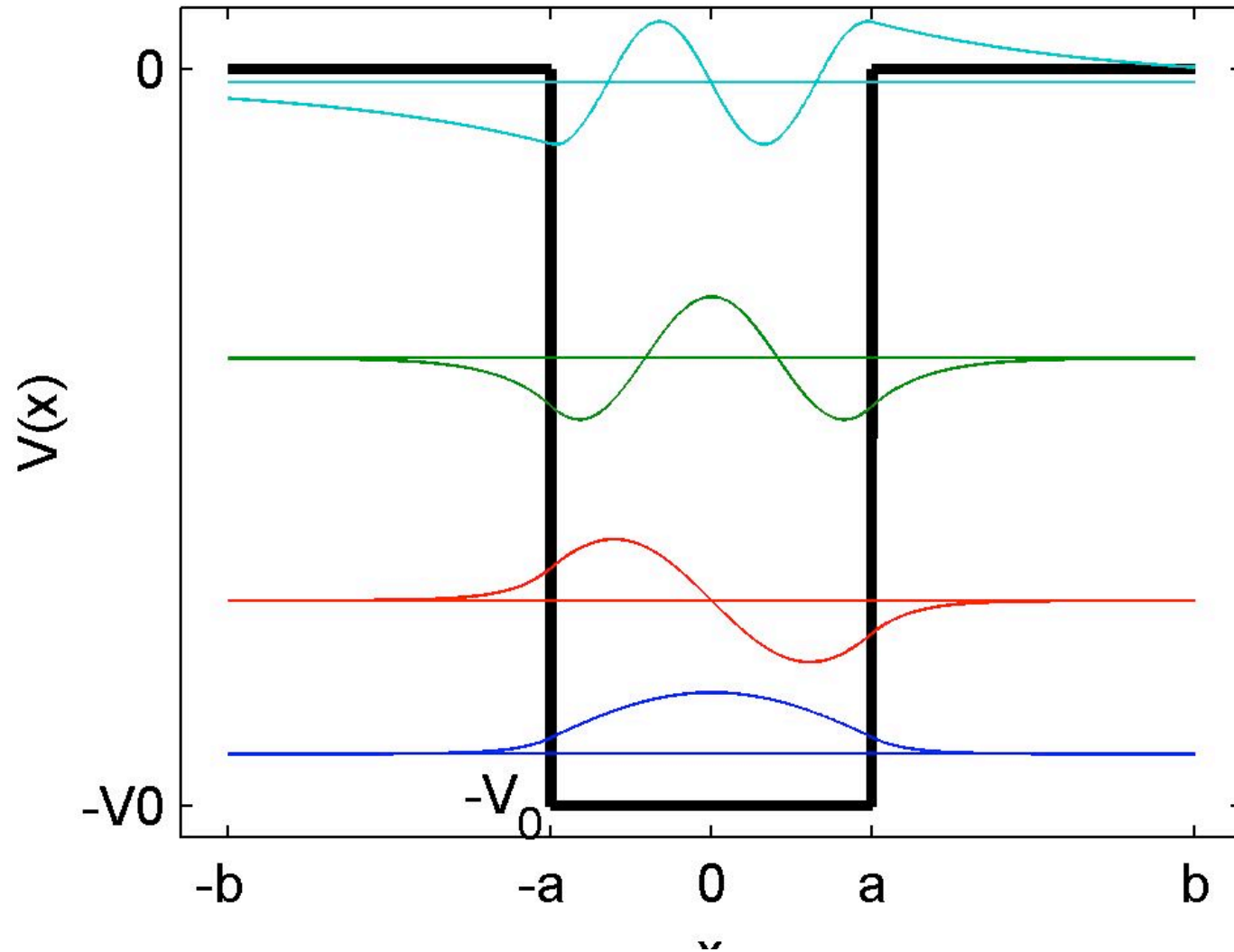
- The number of bound states increases as a function of the strength parameter, and assuming constant mass, the combination

$$V_0a^2$$

is the important parameter, determining the number of bound states, however not in a linear fashion.

Increasing V_0

The square well with bound states



Increasing a

The square well with bound states

